

# Optimal dynamics for quantum information transfer and effective entangling gate through homogeneous quantum wires

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*with*

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**Goal** → information transmission between distant parts

- *Spin chain data-bus*

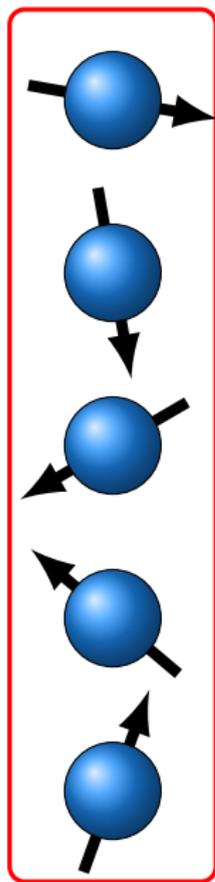
**Tool** → Optimal dynamics

- *induce a coherent (ballistic) dynamics in a dispersive chain*

**Results:**

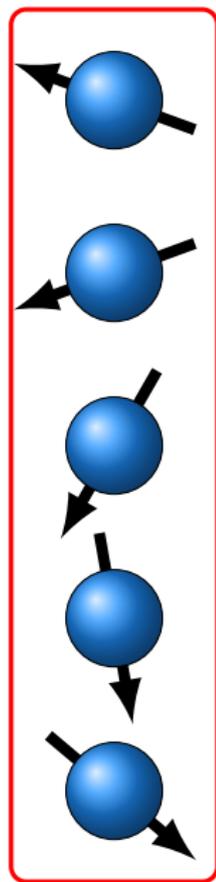
- High quality state transmission
- High quality generation of long-distance entanglement

# Quantum data-bus

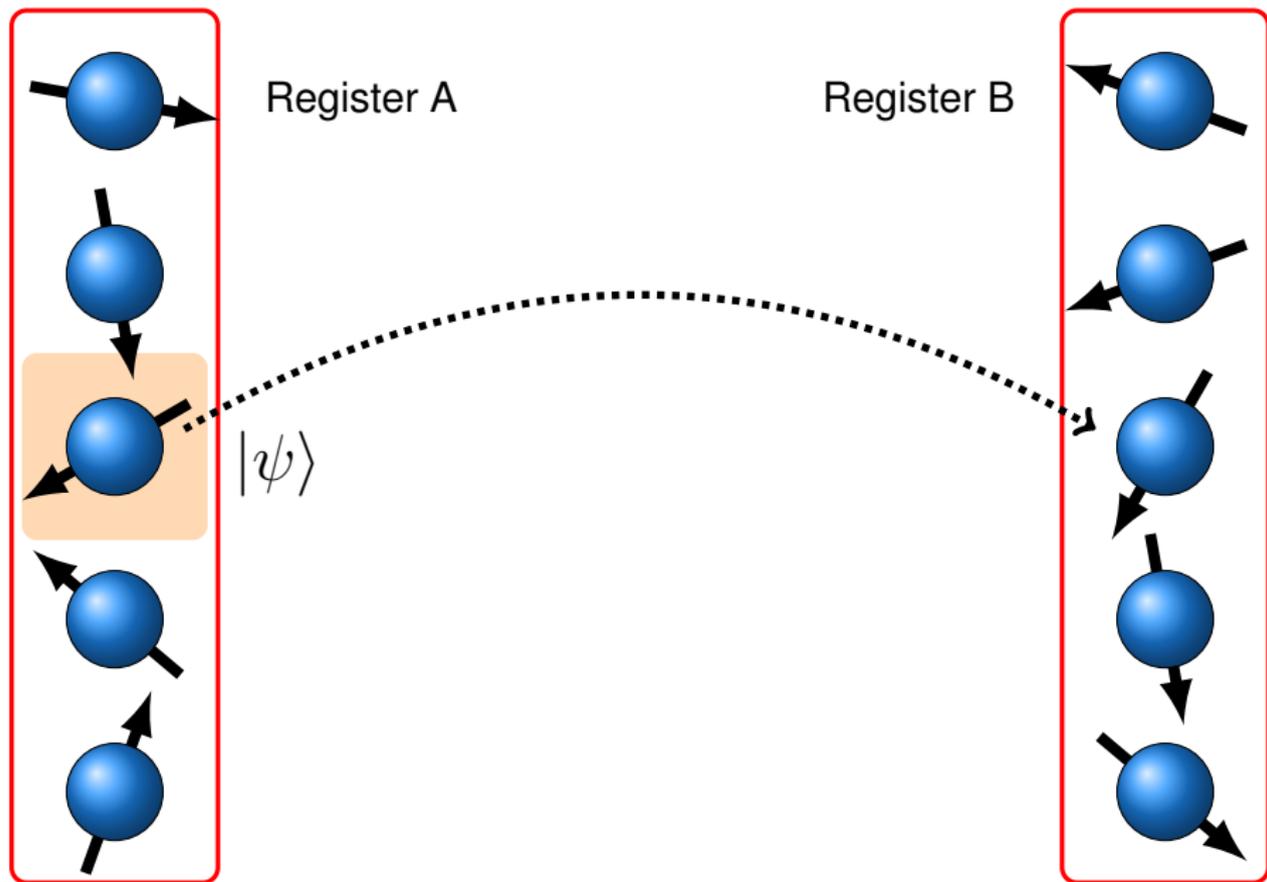


Register A

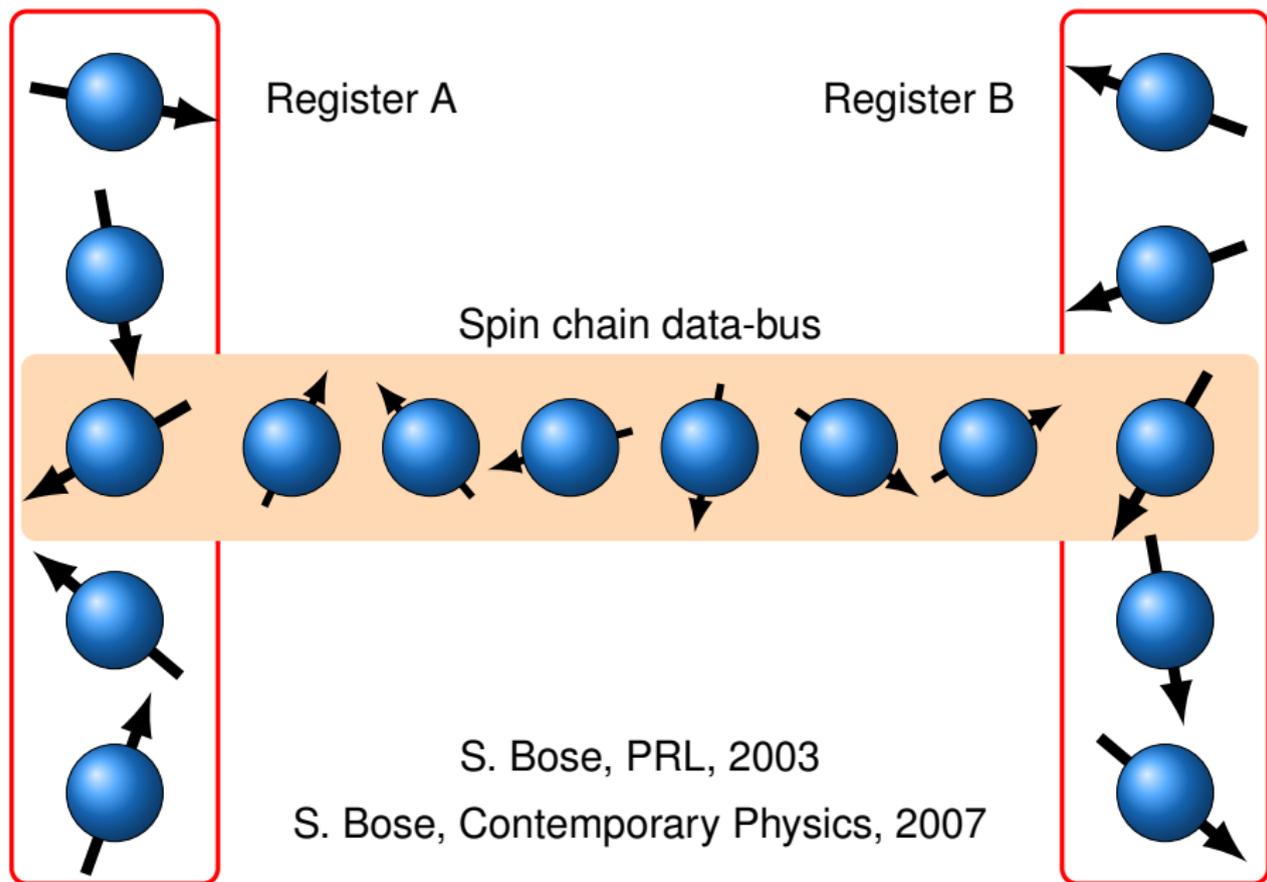
Register B



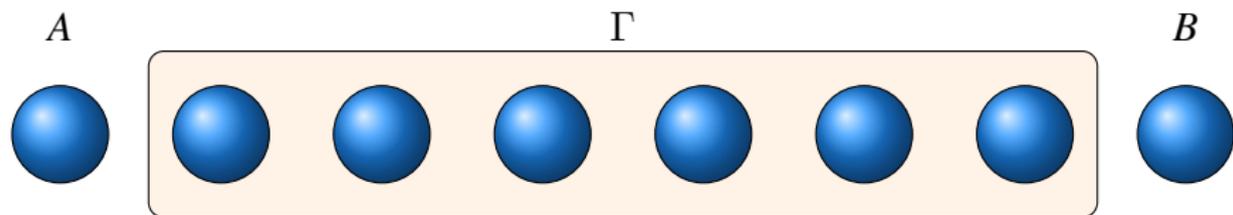
# Quantum data-bus



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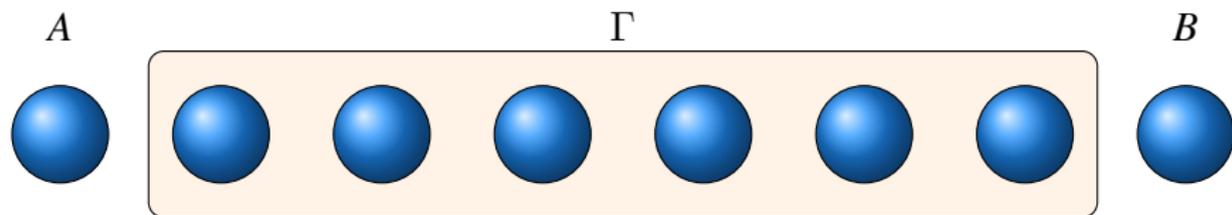


# State transmission



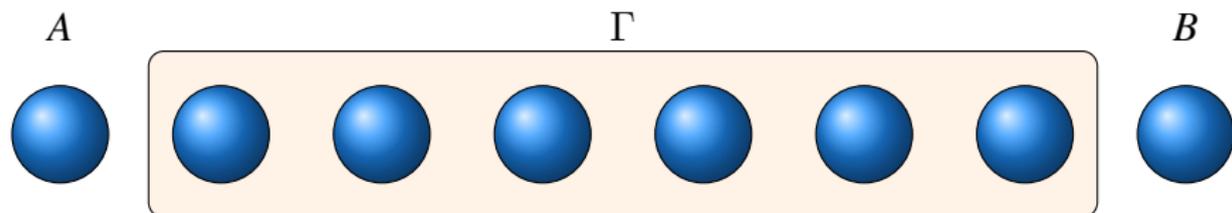
- Out of equilibrium initial state:  $|\Psi_{\text{in}}\rangle = |\psi_A\rangle |\psi_\Gamma\rangle |\psi_B\rangle$

# State transmission



- Out of equilibrium initial state:  $|\Psi_{\text{in}}\rangle = |\psi_A\rangle |\psi_\Gamma\rangle |\psi_B\rangle$
- State of  $B$  at time  $t$ :

$$\rho_B(t) = \text{Tr}_{A\Gamma} \left[ e^{-it\mathcal{H}_{A\Gamma B}} |\Psi_{\text{in}}\rangle \langle \Psi_{\text{in}}| e^{it\mathcal{H}_{A\Gamma B}} \right]$$



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- Measure of the transmission quality: minimum fidelity

$$\mathcal{F}(t) = \min_{|\psi_A\rangle} \langle \psi_A | \rho_B(t) | \psi_A \rangle$$

*Requirement:* quadratic Hamiltonian in terms of *local* creation and annihilation operators

$$\mathcal{H}_{\Lambda\Gamma B} = \sum_{ij} \left( M_{ij} a_i^\dagger a_j + N_{ij} a_i^\dagger a_j^\dagger + \text{h.c.} \right)$$

- Normal modes and dispersion relation  $\omega_k$

$$\mathcal{H}_{\Lambda\Gamma B} = \sum_k \omega_k d_k^\dagger d_k$$

- Density of excitations in the initial state

$$\rho(k) = \langle \Psi_{\text{in}} | d_k^\dagger d_k | \Psi_{\text{in}} \rangle$$

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## Optimal dynamics

Induce a coherent (ballistic) dynamics in a dispersive chain

$\implies$  Make  $\rho(k)$  mostly peaked around the linear zone of  $\omega_k$  through a **minimal** engineering of the interactions

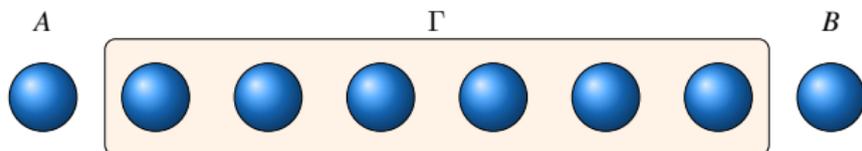
## Optimal dynamics

Induce a coherent (ballistic) dynamics in a dispersive chain

⇒ Make  $\rho(k)$  mostly peaked around the linear zone of  $\omega_k$  through a **minimal** engineering of the interactions

- All relevant excitations (with high  $\rho(k)$ ) have constant group velocity  $v_k = \partial_k \omega_k$
- High quality “wave-packet” transmission
- $\mathcal{F}(t^*) \simeq 1$
- Ballistic transport with arrival time  $t^* \simeq vN$

# Optimal dynamics with the XY model



## XX model

$$\mathcal{H}_\Gamma = J \sum_{n=1}^{N-1} (\sigma_n^x \sigma_{n+1}^x + \sigma_n^y \sigma_{n+1}^y)$$

$$\mathcal{H}_{A\Gamma} = J_0 (\sigma_0^x \sigma_1^x + \sigma_0^y \sigma_1^y)$$

$$\mathcal{H}_{\Gamma B} = J_0 (\sigma_N^x \sigma_{N+1}^x + \sigma_N^y \sigma_{N+1}^y)$$

**Optimal dynamics by setting  $J_0$  to a properly tuned value**

*Banchi, Apollaro, Cuccoli, Vaia, Verrucchi, PRA, 2010*

- Dispersion relation ( $J \equiv 1$ )

$$\omega_k = \cos k$$

- Density of excitations

$$\rho(k) \approx \frac{1}{\Delta^2 + \cot^2 k} \quad \Delta = \frac{J_0^2}{2 - J_0^2}$$

peaked around the inflection point  $k = \pi/2$  with width  $\Delta$

- Normal modes ( $n = 1, \dots, N + 2$ )

$$k_n = \frac{\pi n + \varphi_{k_n}}{N + 3} \quad \varphi_k = 2k - 2 \cot^{-1} \left( \frac{2 - J_0^2}{J_0^2} \cot k \right)$$

*Banchi, Apollaro, Cuccoli, Vaia, Verrucchi, in preparation*

# Optimal dynamics with the XX model

Group velocity of the “wave-packet” close to the inflection point

$$v_k = \partial_k \omega_k \simeq \text{const.} \times \left[ 1 + \left( \frac{2}{NJ_0^6} - \frac{1}{2} \right) k^2 + O(k^4) \right]$$

Optimal values of  $J_0^{\text{opt.}}$ :

$$J_0^{\text{opt.}} \simeq 1.05 N^{-\frac{1}{6}}$$

Non-weak coupling regime and ballistic transport

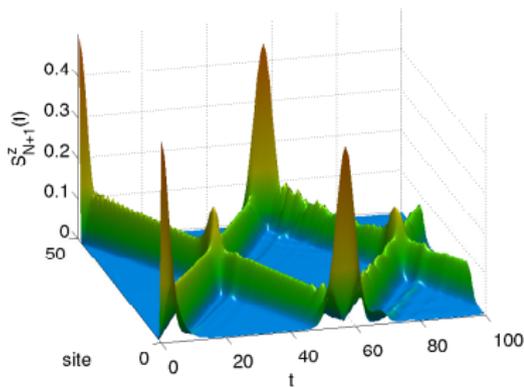
N	25	51	101	251	501	1001	2501	5001
$J_0^{\text{opt.}}$	0.63	0.56	0.49	0.42	0.37	0.33	0.28	0.25

$t^*$	32.5	60.2	112	266	520	1025	2533	5041
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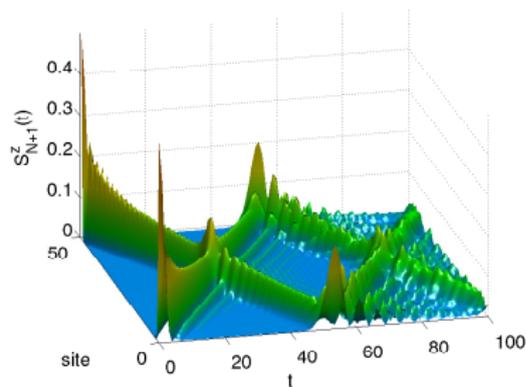


$$|\Psi_{\text{in}}\rangle = |\uparrow\rangle|\Omega_{\Gamma}\rangle|\uparrow\rangle$$

Wave packet propagation displayed by the magnetization  $S^z(x, t)$



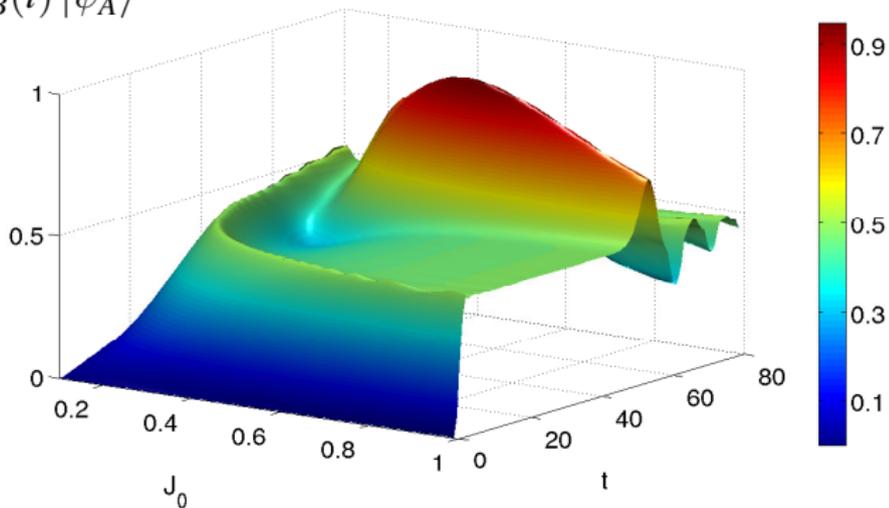
$$J_0 = J_0^{\text{opt.}} \simeq 0.56 J$$



$$J_0 = J$$

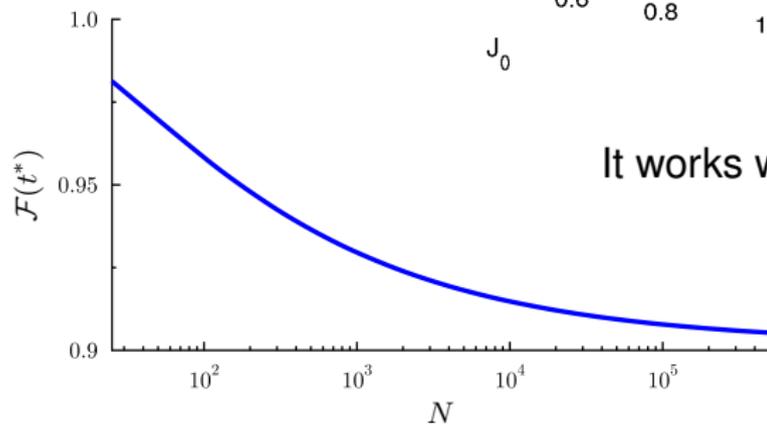
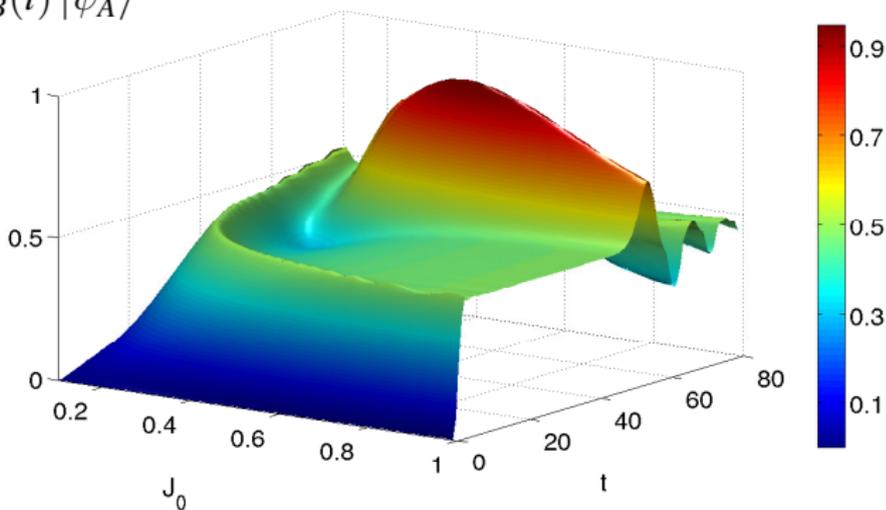
# State transmission quality

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It works with very long channels!

# Entangling quantum gates between distant qubits

$$|\uparrow\rangle_A |\uparrow\rangle_B \longrightarrow e^{i\phi_{\uparrow\uparrow}} |\uparrow\rangle_A |\uparrow\rangle_B$$

$$|\uparrow\rangle_A |\downarrow\rangle_B \longrightarrow e^{i\phi_{\uparrow\downarrow}} |\downarrow\rangle_A |\uparrow\rangle_B$$

$$|\downarrow\rangle_A |\uparrow\rangle_B \longrightarrow e^{i\phi_{\downarrow\uparrow}} |\uparrow\rangle_A |\downarrow\rangle_B$$

$$|\downarrow\rangle_A |\downarrow\rangle_B \longrightarrow e^{i\phi_{\downarrow\downarrow}} |\downarrow\rangle_A |\downarrow\rangle_B$$

Optimal dynamics effectively generates an entangling quantum gate

Fixing

$$|\Psi_{\text{in}}\rangle = (|\uparrow\rangle + |\downarrow\rangle)_A \otimes |\psi\rangle_\Gamma \otimes (|\uparrow\rangle + |\downarrow\rangle)_B$$

$A$  and  $B$  get maximally entangled at  $t^*$

*Banchi, Bayat, Verrucchi, Bose, PRL in print*

$$\mathcal{H} = \sum_{n=0}^N J_n [\sigma_n^x \sigma_{n+1}^x + \sigma_n^y \sigma_{n+1}^y + \Delta \sigma_n^z \sigma_{n+1}^z]$$

where  $J_N = J_0$  and  $J_n = J$  for  $n \neq 0, N$

- Mean field approach for  $\Delta \simeq 0$ : mapping to a quadratic free-fermionic Hamiltonian
- Numerical optimization: finding  $J_0^{\text{opt}}$

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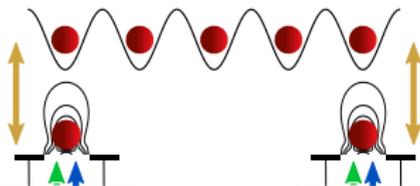
Complex scattering for strong  $\Delta$

- Scattering can favour transmission
- Optimize the initial state of the chain as well

*Bayat, Banchi, Bose, Verrucchi, in preparation*

# Conclusions

- General idea for quadratic Hamiltonians
- High quality information transmission between distant parts
- Experimental feasibility (unmodulated interactions)
  - Cold atoms



*Banchi, Bayat, Verrucchi, Bose, PRL in print*

- Other possible applications in solid state