

Patchy depletion interactions

Antonio Scala

Institute for Complex Systems CNR-ISC
at Univ. di Roma "La Sapienza", Dept. of Physics

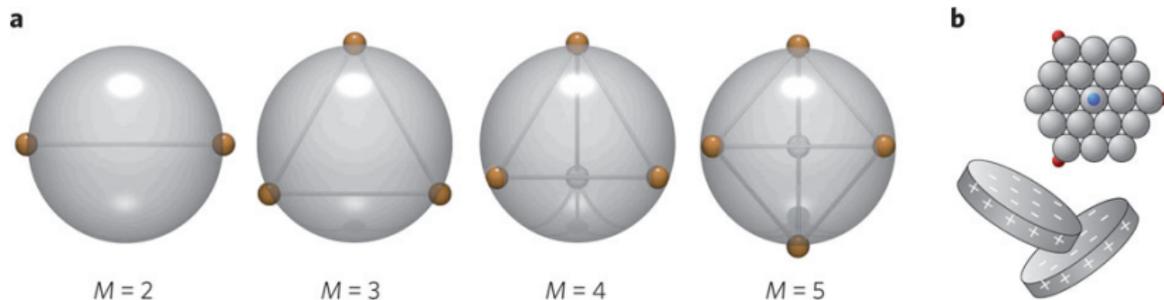
April 4, 2011



Summary

- Shape and Depletion
- Directional depletion in Cut-Spheres
- Key-Lock depletion interaction
- A gallery of shapes
- Conclusions

Which kind of patchyness?



Boltzman envisioned chemical attraction between particles as contact between their sensitive regions, that is, attractive patchy interactions

L. Boltzmann, *Lectures on gas theory*, Dover, New York 1995

A patchy particle is a patterned particle with at least one well-defined patch through which the particle can experience a strongly anisotropic, highly directional interaction with other particles or surfaces

Z. Zhang, S. C. Glotzer, *Nano Lett.* 2004, 4, 1407

The Primitive Model for Water was expanded to to as the Sticky Hard-Sphere model to include an arbitrary number of patches

J. Kolafa, I. Nezbeda, *Mol. Phys.* 1987, 61, 161

Network glasses region in the patchy particle phase diagram

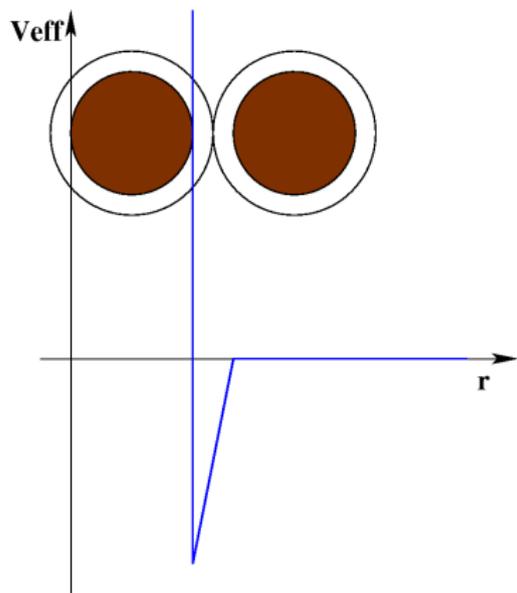
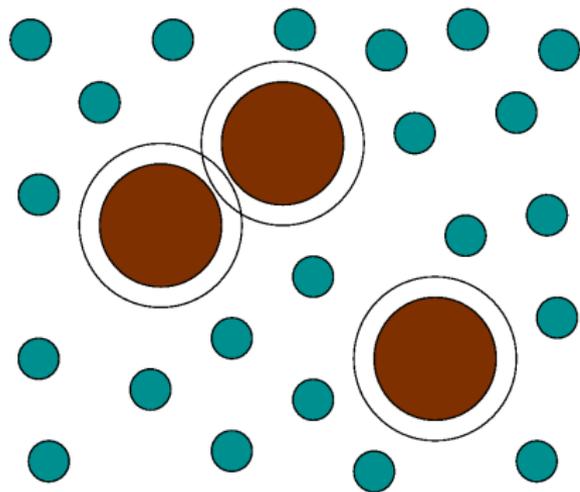
F. Sciortino, *Eur. Phys. J. B* 2008, 64, 505

Empty liquids, Equilibrium Gels

Bianchi et al, *Phys. Rev. Lett.* 2006, 97, 168301

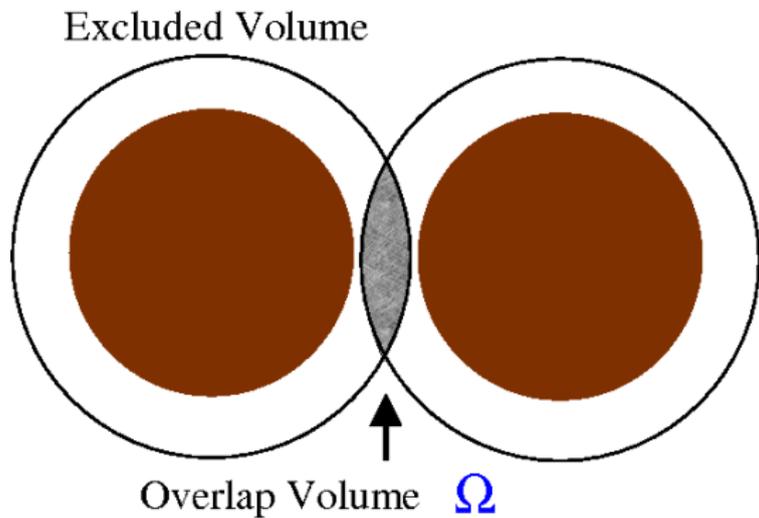
Ruzicka et al, *Nat. Mat.* 2011, 10, 56

Depletion Interactions



- Colloidal Suspensions
- Driven Granulars
- *In-vivo* Proteins ? (*crowded molecular environments*)

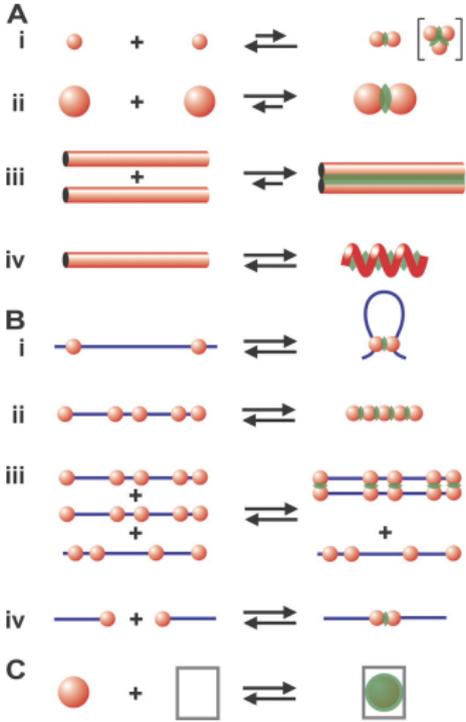
Excluded Volume



Interaction strength $\Delta F \sim f \Omega \sim \phi \Omega$

(free energy per volume $f \sim \phi$ at low density)

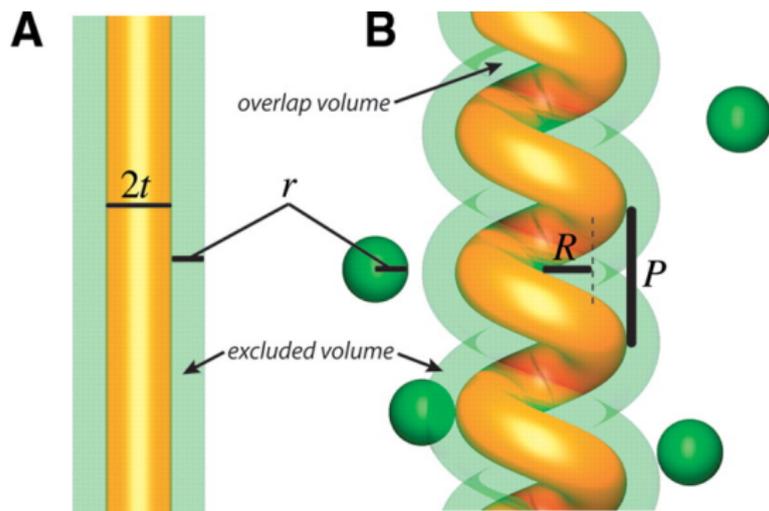
Entropy driven genome organization?



nonspecific entropic forces could also drive biomolecular self-assembly

Daide Marenduzzo et al, *Biophysical Journal* 90, pp. 3712-3721, 2006

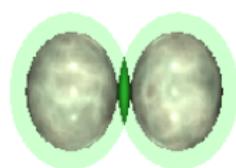
Shape can have dramatic effects



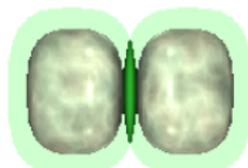
helix forming (Snir & Kamien, Science 307, pp. 1067, 2005)

Flatness enhances interaction

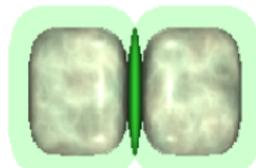
$$\text{SuperEllipsoid: } \left(\frac{x}{a_x}\right)^{n_x} + \left(\frac{y}{a_y}\right)^{n_y} + \left(\frac{z}{a_z}\right)^{n_z} = 1$$



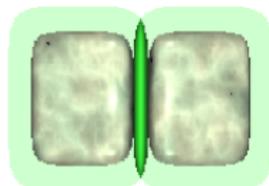
n = 2



n = 3



n = 4

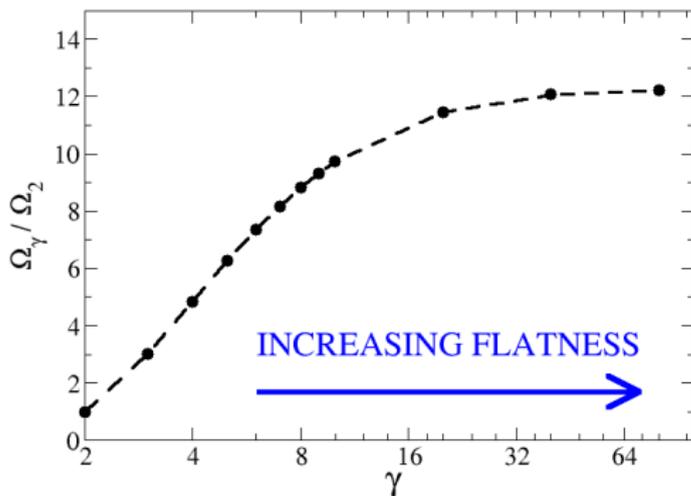


n = 6

- increased strength & induced directionality

Strength vs Flatness

At low depletant density, Ω_γ/Ω_2 estimates the ratio of the interaction strengths



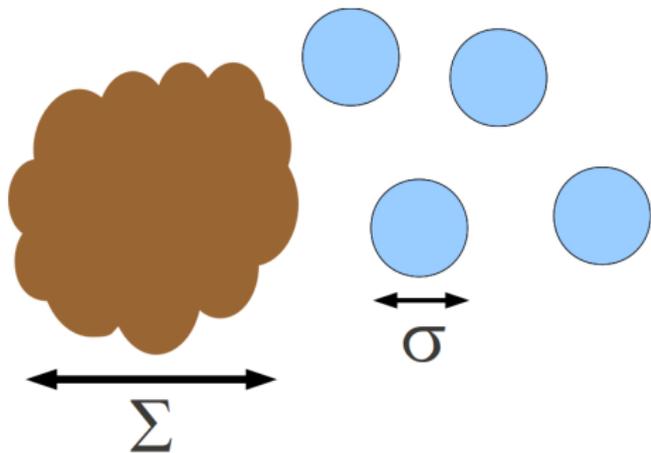
Ω_γ overlap volume of SuperEllipsoids

Ω_2 overlap volume of Spheres

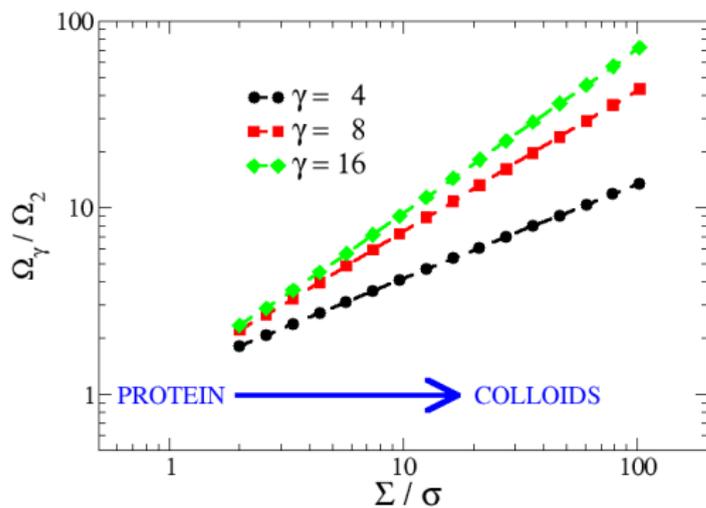
Different scales

protein/nanoparticle limit $\Sigma/\sigma \sim 1$

colloidal limit $\Sigma/\sigma \gg 1$



Strength vs Depletant Size

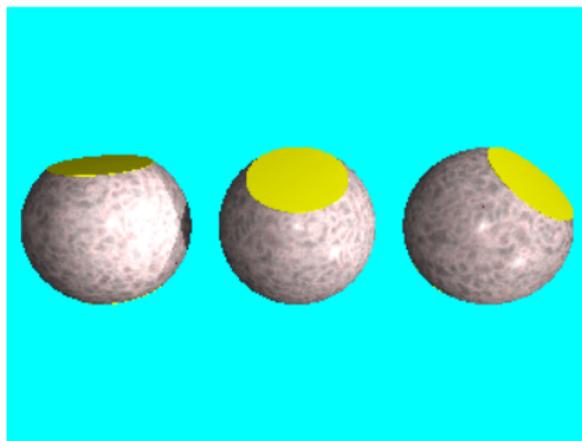
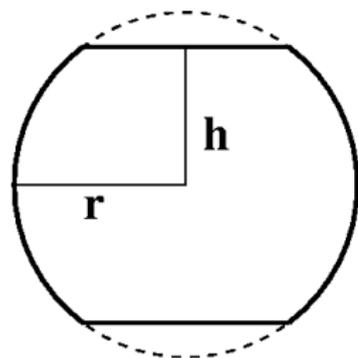


Σ size of depleted particles

σ size of the depletant

Case study 1 : Directional Binding

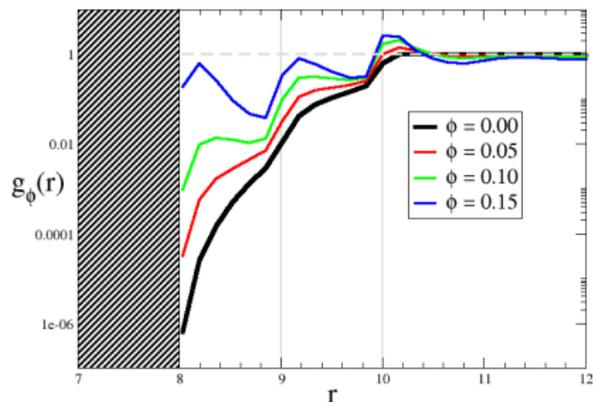
Cut Spheres



We study a template model that highlights the effect shape having:

- either zero curvature (flat faces)
- or constant curvature (spherical sides)

Monte Carlo Simulations

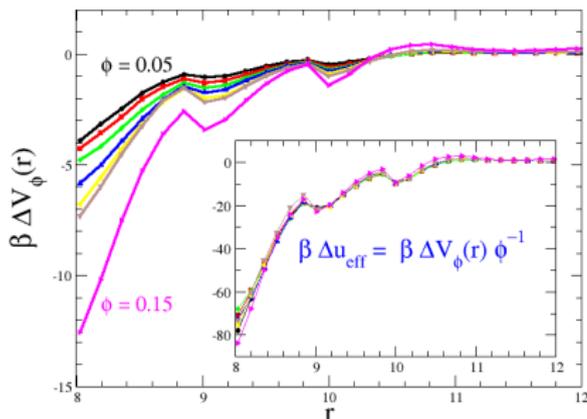


- $\phi = 0.05 \dots 0.15$
- up to 10^4 solvent spheres
- 10 to 20 independent runs
- 10^{11} MC steps per run
- flat $\langle \cos \theta \rangle$
- $\langle \Delta r^2 \rangle \approx L$

Special thanks to Donald Blaak for the CS-CS overlap routine

MC: Scaling of the depletion interaction

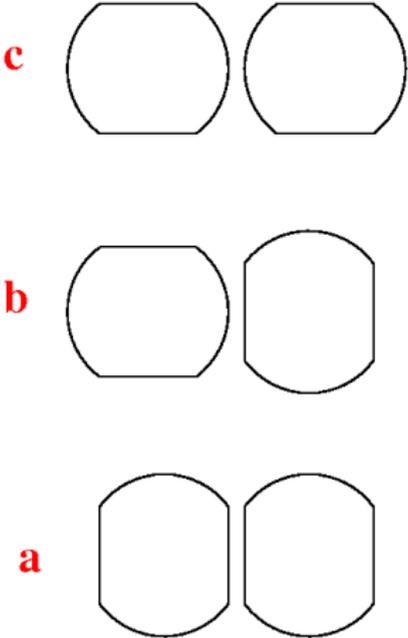
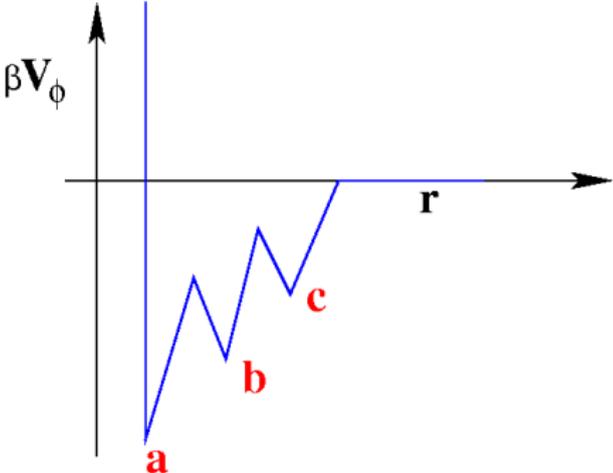
$$g_\phi(r) = g_0(r) \exp[-\beta \Delta V_\phi(r)]$$



For small ϕ , the attractive interaction scales as $\Delta V_\phi \approx \phi \Delta u_{\text{eff}}$

The effective interaction mimics patchy-particles potentials -> toward “empty liquids”

Effective Interaction



Geometry dictates the attractive interaction with three minima among cut spheres

Histogram Method for effective potentials

It is possible to find self-consistently a $\beta V_\phi(r)$ such that sampling moves with probability $\exp[\beta V_\phi(r)]$ the measured radial distribution function is flat
two-body interactions in the dilute limit:

$$g_\phi(r) \equiv_{\text{def}} \exp[-\beta V_\phi(r)]$$

F. Wang and D. P. Landau , Phys. Rev. Lett. 86, p. 2050 , 2001

Bootstrapping the Monte Carlo

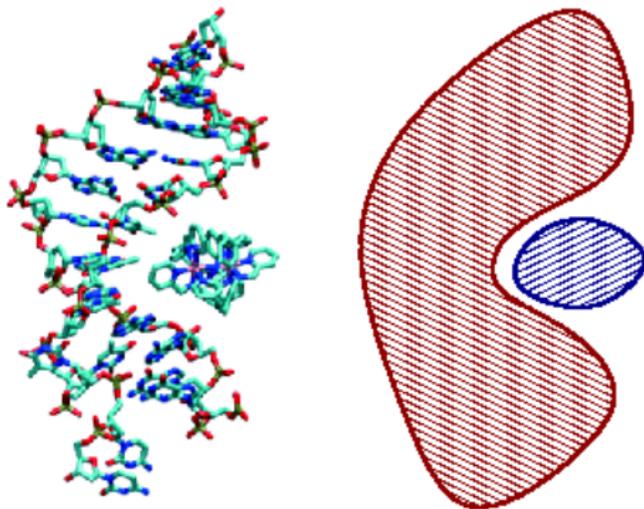
to calculate $\beta \Delta V_{\phi}^{\text{eff}}$ for $\phi = 0 < \phi_1 < \phi_2 < \dots$

- calculate g_0 (simulation of two particles) and set $\beta V_0^{\text{eff}} = -\ln g_0$
- set βV_0^{eff} as initial guess for $\beta V_{\phi_1}^{\text{eff}}$
- calculate g_{ϕ_1}
- set $(\phi_2/\phi_1) \beta V_{\phi_1}^{\text{eff}}$ as initial guess for $\beta V_{\phi_2}^{\text{eff}}$
-

$$\beta \Delta V_{\phi_i}^{\text{eff}} = \beta V_{\phi_i}^{\text{eff}} - \beta V_0^{\text{eff}}$$

Case study 2 : Key-Lock

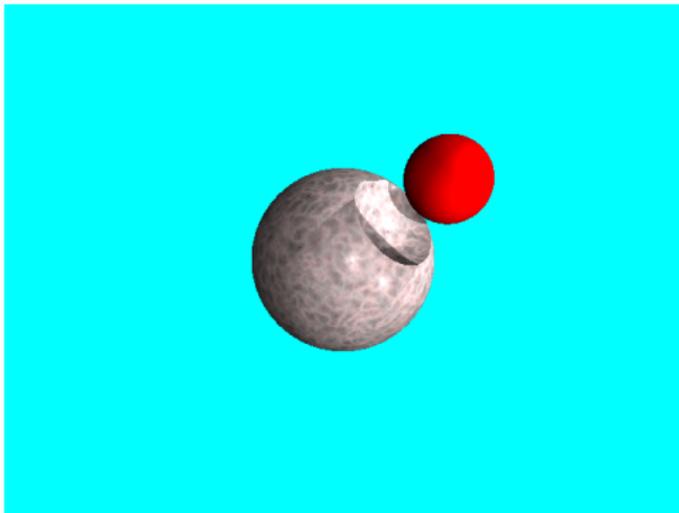
Key-Lock



inspiration: non covalent DNA / metallo-supramolecular cylinder binding

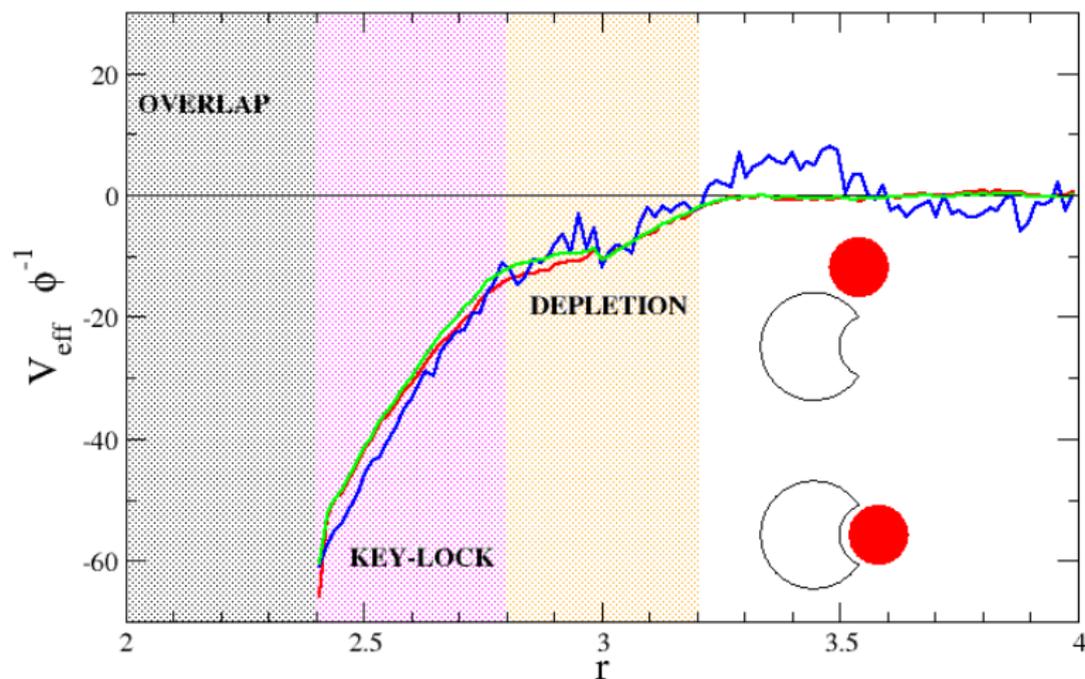
- caveat: real Key-Lock mechanism happen in crowded enviroment and ARE NOT due to depletion
- nevertheless, we test depletion induced key-lock for engineering interactions via shape

Chunk Sphere



MC results

key-lock comes from negative curvature surface exposed to positive curvature surface (even better than flat face - flat face interaction)

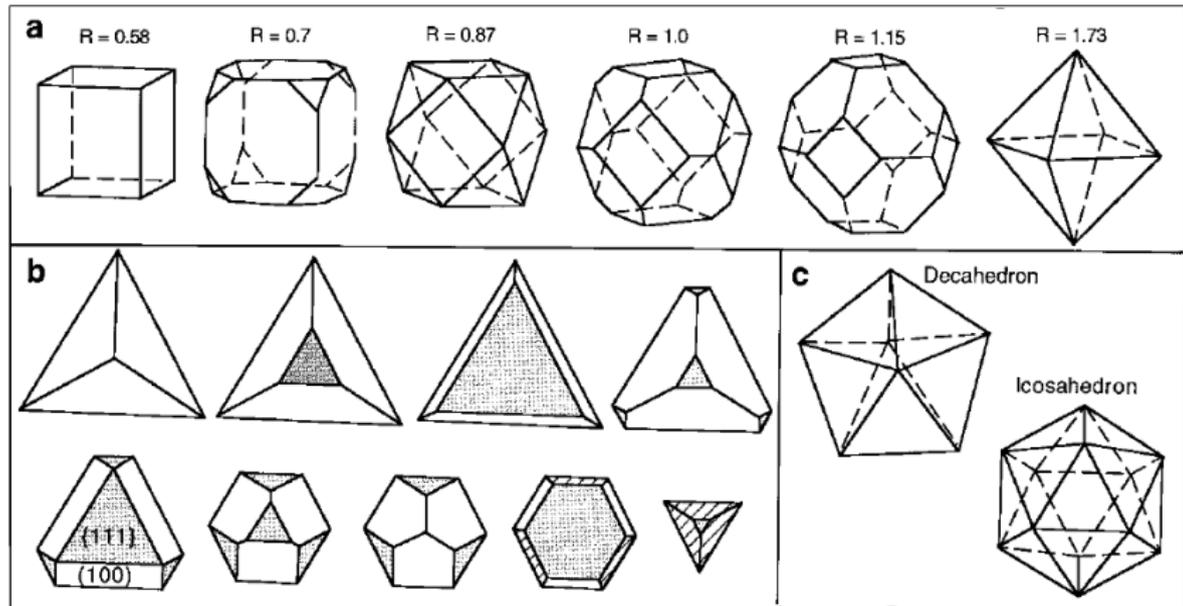


key-lock depletion interaction is ~ 1 order of magnitude bigger

Quest for shaped particles

Nano Crystals

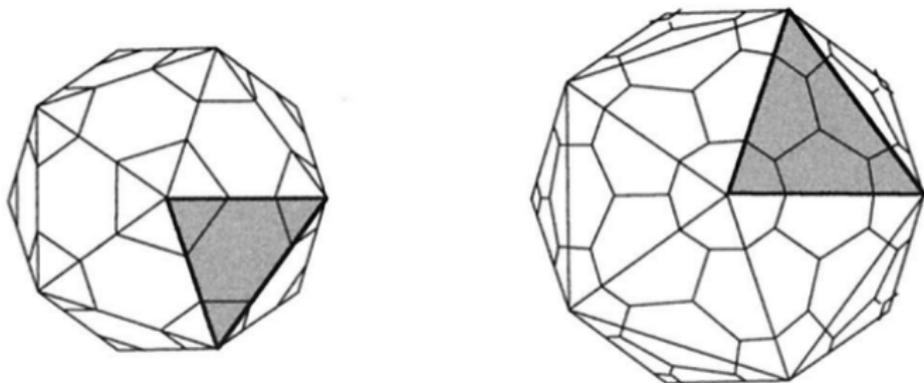
$\approx 10 \text{ nm}$



Zhong Lin Wang, Adv. Mater. 10, pp. 13-30, 1998

Viral Capsids

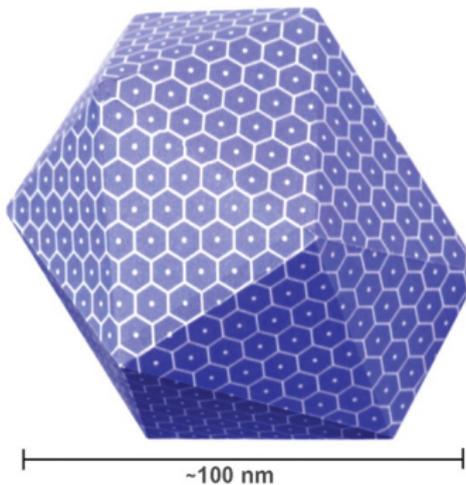
~ 20 nm up to ~ 120 nm



Viral capsids are often monodisperse and highly symmetric nanocontainers

Cellular Microcompartments

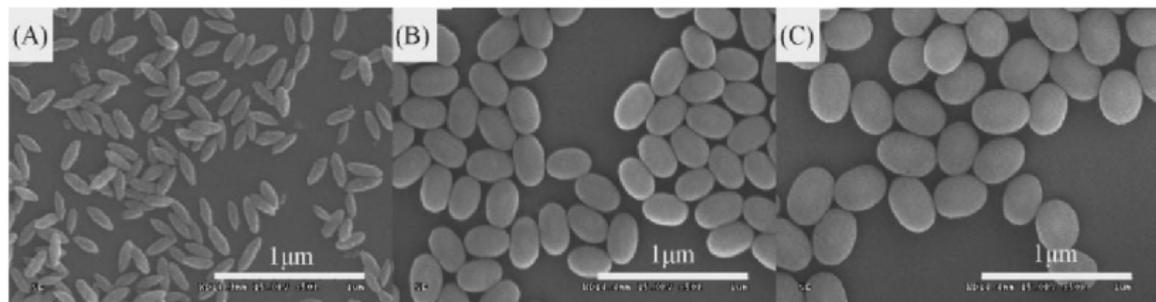
$\sim 100 \text{ nm}$



Shiho Tanaka et al, *Science* 327, pp. 81 - 84, 2010

Colloidal Ellipsoids

$\sim 0.1 \mu\text{m}$ up to $\sim 1 \mu\text{m}$



Tao Ding et al, *Adv. Mater.* 21, pp. 1936–1940, 2009

Conclusions

- Shape allows *directional interactions*
- Shape induces *Key-Lock mechanisms*
- Complex phase diagram (multiple phases)
- Self Assembly
- How to shape particles?
- How to combine the effects of the shape with other interaction?

Conclusions

- Shape allows *directional interactions*
- Shape induces *Key-Lock mechanisms*
- Complex phase diagram (multiple phases)
- Self Assembly
- How to shape particles?
- How to combine the effects of the shape with other interaction?

THANKS !!!

antonio.scala@phys.uniroma1.it